

Solution of the Five-Dimensional Vacuum Einstein Equations in Kaluza–Klein Theory

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In a recent paper Ross obtained the five-dimensional vacuum Einstein equations in Kaluza–Klein theory with energy-momentum tensor equal to zero and solved the equations for a particular case. Here we obtain the complete set of solutions of these equations.

1. INTRODUCTION

In a recent paper Ross (1986) has considered the line element of the form

$$ds^2 = -dt^2 + N(t, \psi) d\psi^2 + R^2(t, \psi) \times [dr^2(1 - kr^2)^{-1} + r^2 d\theta^2 + r^2 \sin^2\theta d\phi^2] \quad (1.1)$$

where r is a dimensionless distance marker and $k = 0, \pm 1$. Ross has shown that for the above metric, the five-dimensional vacuum Einstein equations in Kaluza–Klein theory with A_μ present and a $\partial/\partial\psi$ Killing vector assumed, with energy-momentum tensor $T_{AB} = 0$, reduce to

$$\frac{1}{2} \frac{\ddot{N}}{N} - \frac{1}{4} \left(\frac{\dot{N}}{N} \right)^2 + \frac{3\ddot{R}}{R} = 0 \quad (1.2a)$$

$$-2k + \frac{2R'^2}{N} - 2\dot{R}^2 - R\ddot{R} + \frac{RR''}{N} - \frac{1}{2} \frac{R\dot{R}\dot{N}}{N} - \frac{1}{2} \frac{RR'N'}{N^2} = 0 \quad (1.2b)$$

$$\frac{3R''}{R} - \frac{\ddot{N}}{2} + \frac{\dot{N}^2}{N} - \frac{1}{4} - \frac{3}{2} \frac{R'N'}{RN} - \frac{3}{2} \frac{\dot{R}\dot{N}}{R} = 0 \quad (1.2c)$$

$$\frac{3\dot{R}'}{R} - \frac{3}{2} \frac{\dot{N}R'}{NR} = 0 \quad (1.2d)$$

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where a dot is a time derivative and a prime is a ψ derivative. Equations (1.2a)–(1.2d) were solved by Ross (1986) for the particular case of

$$\begin{aligned} N(\psi, t) &= F(\psi)T(t) \\ R(\psi, t) &= H(\psi)S(t) \end{aligned} \quad (1.3)$$

In the present paper we seek to obtain the complete set of solutions of equations (1.2a)–(1.2d).

2. SOLUTIONS

Integrating equation (1.2d), we obtain

$$N = p(\psi)R'^2 \quad (2.1)$$

where $p(\psi)$ is an arbitrary function of ψ . Using equation (2.1), we can rewrite equation (1.2a) as

$$\frac{\ddot{R}'}{R} = -3 \frac{R'}{R}$$

which on integration gives

$$\ddot{R} = \frac{q(t)}{R^3} \quad (2.2)$$

where $q(t)$ is any arbitrary function of t .

Adding equations (1.2a) and (1.2c) and noting from equation (2.1) that $R' = 0$ leads to $N = 0$, which is impossible, we get

$$\frac{q(t)R'}{R^3} - \dot{R}\dot{R}' - \frac{p'(\psi)}{2p^2(\psi)} = 0$$

Integrating with respect to ψ , we get

$$-\frac{q(t)}{2R^2} - \frac{\dot{R}^2}{2} + \frac{1}{2p(\psi)} = l(t) \quad (2.3)$$

where $l(t)$ is an arbitrary function of t .

Differentiating (2.3) with respect to t and using (2.2), we get

$$-\frac{\dot{q}(t)}{R^2} = 2\dot{l}(t)$$

Since for N not to be zero, R cannot be a function of t only, the above

equation implies

$$\dot{q}(t) = 0 \quad \text{and} \quad \dot{l}(t) = 0$$

i.e.,

$$q(t) = Q, \quad l(t) = L \quad (2.4)$$

where Q and L are integrating constants.

Using equations (2.1)-(2.3), we can rewrite equation (1.2b) in the simple form

$$2L = k \quad (2.5)$$

From equation (2.3) and using (2.4), (2.5), we get

$$R\dot{R} = \left[\left(\frac{1}{p(\psi)} - k \right) R^2 - Q \right]^{1/2}$$

which on integration gives

$$R^2 = \left(\frac{1}{p(\psi)} - k \right) t^2 + 2\mu(\psi)t + \frac{\mu^2(\psi) + Q}{[1/p(\psi)] - k} \quad (2.6)$$

where $\mu(\psi)$ is an arbitrary function of ψ .

Hence, from (2.1),

$$N = p(\psi) \frac{[-(t^2/p^2)p' + 2t\mu' + F'(\psi)]^2}{4[(1/p(\psi) - k)t^2 + 2\mu t + F(\psi)]} \quad (2.7)$$

where

$$F(\psi) = \frac{\mu^2 + Q}{1/p(\psi) - k}$$

It is clear that equations (1.3) represent a special case of equations (2.6) and (2.7) where $\mu = 0$ and $Q = 0$, which is the case considered by Ross (1986).

3. CONCLUSION

We have obtained the complete set of solutions for the five-dimensional vacuum Einstein equations, with zero energy-momentum tensor, given by equations (1.2a)-(1.2d) with metric given by equation (1.1). The complete solutions of equations (1.2a)-(1.2d) are given by equations (2.6) and (2.7).

REFERENCE

Ross, D. K. (1986). *International Journal of Theoretical Physics*, **25**, 663.